

gold wires which were in connexion with the opposite pairs of quadrants of a Thomson's electrometer. The resistance of the film between the needles was calculated by comparing the deflection caused by the difference of potential of the two wires when a current was passing through the film, with that produced by the difference of potential above and below a known resistance placed in the same circuit.

A novel method, the same in principle with the above, was also used to determine the specific resistance of the liquids from which the films were formed. This was deduced from the difference of potential of two platinum wires cemented into a glass tube in which the liquid was contained. As these were at some distance from the electrodes, errors due to polarisation were got rid of. The results of some test experiments made on solutions of sulphuric acid agreed with those of Kohlrausch to within 0·7 per cent.

The authors conclude that their experiments show that the specific resistance of a soap film thicker than  $3\cdot74 \times 10^{-5}$  centims. (the least thickness at which trustworthy observations were made) is independent of the thickness, and is equal to that of the liquid from which it is formed.

They have, therefore, detected no indication of an approach to a thickness equal to the diameter of molecular attraction, and this leads to the deduction that its magnitude must either be less than is supposed by Quincke ( $0\cdot5 \times 10^{-5}$  centims.), or that the mean specific resistance of the surface layer, the thickness of which is equal to that magnitude, does not differ by 17 per cent. from that of the liquid in mass.

They have further found that soap films, even in an enclosed space, may, if the precautions above referred to are not taken, readily lose 23 out of the 57·7 volumes of water contained in every 100 volumes of solution, and their experiments show that this quantity may probably be largely exceeded. They think, therefore, that in all accurate observations on soap films these profound modifications of constitution must either be prevented or measured by a method similar to their own. They criticise from this point of view the observations of Plateau\* and Lüdtge,† and conclude by pointing out the extreme sensitiveness of the electrical method of investigation.

## II. "Molecular Electro-Magnetic Induction." By Professor D. E. HUGHES, F.R.S. Received March 7, 1881.

The induction currents balance which I had the honour of bringing before the notice of the Royal Society‡ showed how extremely sen-

\* "Statique des Liquides," (1873), vol. i, p. 210.

† "Pogg. Ann.," (1870), vol. cxxxvii, p. 620.

‡ "Proc. Roy. Soc.," vol. 29, p. 56.

sitive it was to the slightest molecular change in the composition of any metal or alloy, and it gave strong evidence of a peculiarity in iron and steel which its magnetic properties alone failed to account for. We could with all non-magnetic metals easily obtain a perfect balance of force by an equivalent piece of the same metal, but in the case of iron, steel, and nickel it was with extreme difficulty that I could obtain a near approach to a perfect zero. Two pieces of iron cut off the same bar or wire, possessing the same magnetic moment, never gave identical results; the difficulty was, that notwithstanding each bar or wire could be easily made to produce the same inductive reaction, the time during which this reaction took place varied in each bar; and although I could easily change its balancing power as regards inductive force by a change in the mass of the metal, by heat or magnetism, the zero obtained was never equal to that obtained from copper or silver.

This led me to suppose the existence of a peculiarity in magnetic metals which could not be accounted for except upon the hypothesis that there was a cause, then unknown, to produce the invariable effect.

Finding that it would be impossible to arrive at the true cause without some new method of investigation, which should allow me to observe the effects from an electrical circuit, whose active portion should be the iron wire itself, I constructed an apparatus or electro-magnetic induction balance, consisting of a single coil reacting upon an iron wire in its axis, and perpendicular to the coil itself; by this means the iron or other wire itself became a primary or secondary, according as the current passed through the coil or wire. Now, with this apparatus we could induce secondary currents upon the wire or coil, if the coil was at any angle, except when the wire was absolutely perpendicular; in this state we could only obtain a current from some disturbing cause, and the current so obtained was no longer secondary but tertiary.

The whole apparatus, however, is more complicated than the general idea given above, as it was requisite not only to produce effects but to be able to appreciate the direction of the electrical current obtained, and have comparative measures of their value. In order to fully understand the mode of experiments, as well as the results obtained, I will first describe the apparatus employed.

The electro-magnetic induction balance consists of—1st, an instrument for producing the new induction current; 2nd, sonometer or balancing coils; 3rd, rheotome and battery; 4th, telephone.

The essential portion of this new balance is that wherein a coil is so arranged that a wire of iron or copper can pass freely through, forming its axis. The iron or copper wire rests upon two supports 20 centims. apart; at one of these the wire is firmly clamped by two binding screws; the opposite end of the wire turns freely on its sup-

port, the wire being 22 centims. long, having 2 centims. projection beyond its support, in order to fasten upon it a key or arm which shall serve as a pointer upon a circle giving the degrees of torsion which the wire receives from turning this pointer. A binding screw allows us to fasten the pointer at any degree, and thus preserve the required stress the time required.

The exterior diameter of the coil is  $5\frac{1}{2}$  centims., having an interior vacant circular space of  $3\frac{1}{2}$  centims., its width is 2 centims.; upon this is wound 200 metres of No. 32 silk-covered copper wire. This coil is fastened to a small board so arranged that it can be turned through any desired angle in relation to the iron wire which passes through its centre, and it can also be moved to any portion of the 20 centims. of wire, in order that different portions of the same wire may be tested for a similar stress.

The whole of this instrument, as far as possible, should be constructed of wood, in order to avoid, as far as possible, all disturbing inductive influences of the coil upon them.

The iron wire at its fixed end is joined or makes contact with a copper wire, which returns to the front part of the dial under its board and parallel to its coil, thus forming a loop, the free end of the iron wire is joined to one pole of the battery, the copper wire under the board is joined to the rheotome and thence to the battery.

The coil is joined to the telephone; but, as in every instance we can either pass the battery current through the wire, listening to its inductive effects upon the wire, or the reverse of this; I prefer, generally, in order to have no voltaic current passing through the wire, to join the iron wire and its loop direct to the telephone, passing the voltaic current through the coil.

In order to balance, measure, and know the direction of the new induction currents by means of a switching key, the sonometer\* I described to the Royal Society is brought into the circuit. The two exterior coils of the sonometer are then in the circuit of the battery, and of the coil upon the board containing the iron wire or stress bridge. The interior or movable coil of the sonometer is then in the circuit of the iron wire and telephone. Instead of the sonometer constructed as described in my paper to the Royal Society, I prefer to use one formed upon a principle I described in "Comptes Rendus," December 30, 1878. This consists of two coils only, one of which is smaller and turns freely in the centre of the outside coil. The exterior coil being stationary, the centre coil turns upon an axle by means of a long (20 centims.) arm or pointer, the point of which moves over a graduated arc or circle. Whenever the axis of the interior coil is perpendicular to the exterior coil, no induction takes place, and we

\* "Proc. Roy. Soc." vol. 29, p. 65.

have a perfect zero : by turning the interior coil through any degree we have a current proportional to this angle, and in the direction in which it is turned. As this instrument obeys all the well-known laws for galvanometers, the readings and evaluations are easy and rapid.

If the coil upon the stress bridge is perpendicular to the iron wire, and if the sonometer coil is at zero, no currents or sounds in the telephone will be perceived, but the slightest current in the iron wire produced by torsion will at once be heard ; and by moving the sonometer coil in a direction corresponding to the current, a new zero will be obtained, which will not only balance the force of the new current but indicate its value. A perfect zero, however, will not be obtained with the powerful currents obtained by the torsion of 2 milims. diameter iron wire, we then require special arrangements of the sonometer which are too complicated to describe here.

The rheotome is a clockwork, having a rapid revolving wheel which gives interruptions of currents in fixed cadences in order to have equal intervals of sound and silence. I employ four bichromate cells or eight Daniell's elements, and they are joined through this rheotome to the coil on the stress bridge, as I have already described.

The magnetic properties of iron, steel, nickel, and cobalt, have been so searchingly investigated by ancient as well as by modern scientific authors, that there seems little left to be known as regards its molar magnetism. I use the word molar here simply to distinguish or separate the idea of a magnetic bar of iron or steel magnetised longitudinally or transversely from the polarised molecules which are supposed to produce its external magnetic effects.

Molar magnetism, whilst having the power of inducing an electric current in an adjacent wire, provided that either has motion, or the magnet a change in its magnetic force, as shown by Faraday in 1832, has no power of inducing an electric current upon itself or its own molar constituent, either by motion or change of its magnetic moment. Molecular magnetism (the results of which, I believe, I have been the first to obtain) has no, or a very feeble, power of inducing either magnetism or an electric current in an adjacent wire, but it possesses the remarkable power of strongly reacting upon its own molar wire, inducing (comparatively with its length) powerful electric currents, in a circuit of which this forms a part.

In some cases, as will be shown, we may have both cases existing in the same wire ; this occurs when the wire is under the influence of stress, either external or internal ; it would have been most difficult to separate these two, as it was in my experiments with the induction balance, without the aid of my new method.

Ampère's theory supposes a molecular magnetism or polarity, and that molar magnetism would be produced when the molecular mag-

netism became symmetrical ; and his theory, I believe, is fully capable of explaining the effects I have obtained, if we admit that we can rotate the paths of the polarised molecules by an elastic torsion.

Matteucci made use of an inducing and secondary coil in the year 1847,\* by means of which he observed that mechanical strains increased or depressed the magnetism of a bar inside this coil.

Wertheim published in the "Comptes Rendus," 1852,† some results similar to Matteucci ; but in the "Annales de Chimie et de Physique," 1857,‡ he published a long series of most remarkable experiments, in which he clearly proves the influence of torsion upon the increment or decrement of a magnetical wire.

Vilari showed§ increase or diminution of magnetism by longitudinal pull according as the magnetising force is less or greater than a certain critical value.

Wiedemann,|| in his remarkable work "Galvanismus," says that an iron wire through which an electric current is flowing becomes magnetised by twisting the wire. This effect I have repeated, but found the effects very weak, no doubt due to the weak battery I use, viz., four quart bichromate cells.

Sir W. Thomson shows clearly in his remarkable paper "Effects of Stress on the Magnetisation of Iron, Nickel, and Cobalt,"¶ the critical value of the magnetisation of these metals under varying stress, and also explains the longitudinal magnetism produced by Wiedemann as due to the outside molar twist of the wire, making the current pass as in a spiral round a fixed centre. Sir William Thomson also shows clearly the effects of longitudinal as well as transversal strain, both as regards its molar magnetism and its electric conductivity.

My own researches convince me that we have in molecular magnetism a distinct and separate form of magnetism from that when we develop, or render evident, longitudinal or transversal magnetism, which I have before defined as molar.

Molecular magnetism is developed by any slight strain or twist other than longitudinal, and it is only developed by an elastic twist; for, however much we may twist a wire, provided that its fibres are not separated, we shall only have the result due to the reaction of its remaining elasticity.

If we place an iron wire, say 20 centims. long, 1 millim. diameter, in the axis of the coil of the electro-magnetic balance, and if this wire is joined, as described, to the telephone, we find that on passing a electric

\* "Compt. Rend.," t. xxiv, p. 301, 1847.

† "Compt. Rend.," t. xxv, p. 702, 1852.

‡ "Ann. de Chim. et de Phys.," (3), t. 1, p. 385, 1857.

§ "Poggendorff's Annalen," 1868.

|| Wiedemann. "Galvanismus," p. 447.

¶ "Phil. Trans.," 1879, p. 55.

current through the inducing coil no current is perceptible upon the iron wire; but if we give a very slight twist to this wire at its free end—one-eighteenth of a turn, or  $20^\circ$ —we at once hear, clear and comparatively loud, the currents passing the coil; and although we only gave a slight elastic twist of  $20^\circ$  of a whole turn, and this spread over 20 centims. in length, making an extremely slight molar spiral, yet the effects are more powerful than if, using a wire free from stress, we turned the whole coil  $40^\circ$ . The current obtained when we turn the coil, as just mentioned, is secondary, and with the coil at any angle any current produced by its action, either on a copper, silver, iron, or steel wire; in fact, it is simply Faraday's discovery; but the current from an elastic twist is no longer secondary under the same conditions, but tertiary, as I shall demonstrate later on. The current passing through the coil cannot induce a current upon a wire perpendicular to itself, but the molecules of the outside of the wire, being under a greater elastic stress than the wire itself, they are no longer perpendicular to the centre of the wire, and consequently they react upon this wire as separate magnets would upon an adjacent wire. It might here be readily supposed that a wire having several twists, or a fixed molar twist of a given amount, would produce similar effects. It, however, does not, for in most cases the current obtained from the molar twists are in a contrary direction to that of the elastic torsion. Thus, if I place an iron wire under a right-handed elastic twist of  $20^\circ$ , I find a positive current of  $50^\circ$  sonometer; but if I continue this twist so that the index makes one or several entire revolutions, thus giving a permanent molar twist of several turns, I find upon leaving the index free from any elastic torsion, that I have a permanent current of  $10^\circ$ , but it is no longer positive, but negative, requiring that we should give an elastic torsion in the previous direction, in order to produce a positive current. Here a permanent elastic torsion of the molecules is set up in the contrary direction to its molar twist, and we have a negative current, overpowering any positive current which should have been due to the twisted wire.

The following table shows the influence of a permanent twist, and that the current obtained when the wire was freed from its elastic torsion was in opposition to that which should have been produced by the permanent twist. Thus, a well-softened iron wire, 1 millim. in diameter, giving  $60^\circ$  positive current for a right-handed elastic torsion of  $20^\circ$ , gave after  $1^\circ\cdot80$  permanent torsion a negative current of  $10^\circ$ .

1	complete permanent torsion (right-handed)	negative	....	10
2	"	"	"	.... 15
3	"	"	"	.... 15
4	"	"	"	.... 16
5	"	"	"	.... 12

6	complete permanent torsion (right-handed)	negative	....	10
7	"	"	"	.... 5
8	"	"	"	.... 4
9	"	"	"	.... 3
10	"	"	"	.... 3

At this point the fibres of a soft wire commence to separate, and we have no longer a complete single wire, but a helix of separate wires upon a central structure.

If now, instead of passing the current through the coil, I pass it through the wire, and place the telephone upon the coil circuit, I find that I obtain equally as strong tertiary currents upon the coil as in the previous case, although in the first case there was produced longitudinal electro-magnetism in the perpendicular wire by the action of the coil, but in the latter case none or the most feeble electro-magnetism was produced, yet in these two distinct cases we have a powerful current produced not only upon its own wire, but upon the coil, thus proving that the effects are equally produced both on the wire and coil.

If we desire, however, in these reversible effects to produce in both cases the same electromotive force, we must remember that the tertiary current when reacting upon its own short wire produces a current of great quantity, the coil one of comparative higher intensity. We can, however, easily convert the great quantity of the wire into one of higher tension by passing it through the primary of a small induction coil whose resistance is not greater than one ohm. We can then join our telephone, which may be then one of a high resistance, to the secondary of this induction coil, and by this means, and without changing the resistance of the telephone, receive the same amount of force, either from the iron wire or the coil.

Finding that iron, steel, and all magnetic metals produce a current by a slight twist, if now we replace this wire by one of copper or non-magnetic metals we have no current whatever by an elastic twist, and no effects, except when the wire itself is twisted spirally in helix, and whatever current we may obtain from copper, &c., no matter if from its being in spiral or from not being perpendicular to the axis of the coils, the currents obtained will be invariably secondary and not tertiary. If we replace the copper by an iron wire, and give it a certain fixed torsion, not passing its limit of elasticity, we find that no increase or decrease takes place by long action or time of being under strain. Thus a wire which gave a sonometric force of  $50^\circ$  at the first observation, remained perfectly constant for several days until it was again brought to zero by taking off the strain it had received. Thus we may consider that as long as the wire preserves its elasticity, exactly in the same ratio will it preserve the molecular character of its magnetism.

It is not necessary to use a wire to produce these effects ; still more powerful currents are generated in bars, ribbons, or sheets of iron ; thus, no matter what external form the iron may possess, it still produces all the effects I have described.

It requires a great many permanent twists in a wire to be able to see any effect from these twists, but if we give to a wire, 1 millim. diameter, forty whole turns (or until its fibres become separated) we find some new effects ; we find a small current of  $10^\circ$  in the same direction as its molar twist, and on giving a slight twist ( $20^\circ$ ) the sonometric value of the sound obtained is  $80^\circ$ , instead of  $50^\circ$ , the real value of a similar untwisted wire ; but the explanation will be found by twisting the wire in a contrary direction to its molar twist. We can now approach the zero but never produce a current in the contrary direction, owing to the fact that by the spiral direction, due to the fibrous molar turns, the neutral position of its molecules is no longer parallel with its wire, but parallel with its molar twist, consequently an elastic strain in the latter case can only bring the molecules parallel with its wire, producing no current, and in the first case the angle at which the reaction takes place is greater than before, consequently the increased value of its current.

The measurements of electric force mentioned in this paper are all sonometric on an arbitrary scale. Their absolute value has not yet been obtained, as we do not, at our present stage, require any except comparative measures.\* Thus, if each wire is of 1 millim. diameter and 20 centims. long, all render the same stress in the axis of its coil. I find that the following are the sonometric degrees of value :—

	<sup>c</sup>	Tertiary current.
Soft iron . . . . .	60	
Hard drum iron . . . . .	50	"
Soft steel . . . . .	45	"
Hard tempered steel . . . . .	10	"
Copper, silver, &c. . . . .	-0	
Copper helix, 1 centim. diameter, 20 turns in 20 centims. . . . .	45	Secondary currents.
Iron, spiral, ditto . . . . .	45	"
Steel . . . . .	45	"

The tertiary current increases with the diameter of the wire, in a proportion which has not yet been determined ; thus, an ordinary hard iron wire of 1 millim. diameter giving  $50^\circ$ , one of 2 millims. diameter gave  $100^\circ$  ; and the maximum of force obtained by any degree of torsion is at or near the limit of elasticity, since as soon as we pass this point, producing a permanent twist, the current decreases, as I have already shown in the case of a permanent twist. Thus, the

\*  $50^\circ$  sonometer has the same electromotive force as 0'10 of a Daniell battery.

critical point of 1 millim. hard iron wire was  $20^\circ$  of torsion, but in hard steel it was  $45^\circ$ .

Longitudinal strains do not produce any current whatever, but a very slight twist to a wire, under a longitudinal strain, produces its maximum effects: thus,  $20^\circ$  of torsion being the critical point of iron wire, the same wire, under longitudinal strain, required but from  $10^\circ$  to  $15^\circ$ . It is very difficult, however, to produce a perfect longitudinal strain alone. I have, therefore, only been able to try the effect of longitudinal strain on fine wires, not larger than 1 millim. in diameter, but as in all cases, no effect whatever was produced by longitudinal strain alone, I believe none will be found if the wire be absolutely free from torsion. The molecules in a longitudinal strain are equally under an elastic strain as in torsion, but the path of their motion is now parallel with its wire, or the zero of electric inductive effect, but the longitudinal and transverse strains of which the compound strain is composed, react upon each other, producing the increased effect due to the compound strain.

The sonometer is not only useful for showing the direction of the current and measuring it by the zero method, but it also shows at once if the current measured is secondary or tertiary. If the current is secondary its period of action coincides with that of the sonometer, and a perfect balance, or zero of sound, is at once obtained, and its value in sonometric degrees given, but if the current is tertiary, no zero is possible, and if the value of the tertiary is  $60^\circ$ , we find  $60^\circ$  the nearest approach to zero possible. But by the aid of separate induction coils to convert the secondary into a tertiary, a perfect zero can be obtained if the time of action and its force correspond to that which we wish to measure.

If I place a copper wire in the balance and turn the coils at an angle of  $45^\circ$ , I obtain a current for which the sonometer gives a perfect zero at  $50^\circ$ , proving, as already said, that it is secondary. If I now replace the copper by an iron wire, the coil remaining at  $45^\circ$ , I have again exactly the same value for the iron as copper, viz.,  $50^\circ$ , and in both cases secondary. Now, it is evident that in the case of the iron wire there was produced at each passage of the current a strong electro-magnet, but this longitudinal magnetism did not either change the character of the current or its value in force.

A most beautiful demonstration of the fact that longitudinal magnetism produces no current, but that molecular magnetism can act equally as well, no matter the direction of the longitudinal magnetism, consists in forming an iron wire in a loop, or taking two parallel but separate wires, joined electrically at their fixed ends, the free ends being each connected with the circuit, so that the current generated must pass up one wire and down the adjacent one. On testing this loop, and if there are no internal strains, complete silence or absence

of current will be found. Now, giving a slight torsion to one of these wires in a given direction, we find, say,  $50^\circ$  positive; twisting the parallel wire in a similar direction produces a perfect zero, thus, the current of the second must have balanced the positive of the first. If, instead of twisting it in similar directions, we twist it in the contrary direction, the sounds are increased in value from  $50^\circ$  positive to  $100^\circ$  positive, showing, in this latter case, not only a twofold increase of force, but that the currents in the iron wires travelled up one wire and down the other, notwithstanding that both were strongly magnetic by the influence of the coil in one direction, and this experiment also proves that its molar magnetism had no effect, as the currents are equally strong in both directions, and both wires can double or efface the currents produced in each. If, instead of two wires we take four, we can produce a zero, or a current of  $200^\circ$ , and with twenty wires we have a force of  $1,000^\circ$ , or an electromotive force of two volts. We have here a means of multiplying the effects by giving an elastic torsion to each separate wire, and joining them electrically in tension. If loops are formed of one iron and one copper wire, we can obtain both currents from the iron wire, positive and negative, but none from the copper, its *rôle* is simply that of a conductor upon which torsion has no effect.

I have already mentioned that internal strains will give out tertiary currents, without any external elastic strain being put on. In the case of iron wire, these disappear by a few twists in both directions, but in flat bars or forged iron, they are more permanent; evidently, portions of these bars have an elastic strain, whilst other portions are free, for I find a difference at every inch tested: the instrument, however, is so admirably sensitive, and able to point out not only the strain, but its direction, that I have no doubt its application to large forged pieces, such as shafts or cannon, would bring out most interesting results, besides its practical utility; great care is therefore necessary in these experiments that we have a wire free from internal strains, or that we know their value.

Magnetising the iron wire by a large steel permanent magnet has no effect whatever. A hard steel wire thus placed becomes strongly magnetic, but no current is generated, nor has it any influence upon the results obtained from molecular movement, as in elastic torsion. A flat wide iron or steel bar shows this better than iron wire, as we can here produce transversal, instead of longitudinal, but neither shows any trace of the currents produced by molecular magnetism. I have made many experiments with wires and bars thus magnetised, but as the effect in every case was negative when freed from experimental errors, I will not mention them; but there is one very interesting proof which the instrument gives, that longitudinal magnetism first passes through its molecular condition before and during the discharge,

or recombination of its magnetism. For this purpose, using no battery, I join the rheotome and telephone to the coil, the wire having no exterior circuit. If I strongly magnetise the two ends of the wire, I find by rapidly moving the coil, that there is a Faradaic induction of  $50^\circ$  at both poles, but very little or none at the centre of the wire; now fixing the coil at the centre or neutral point of the wire, and listening intently, no sounds are heard, but the instant I give a slight elastic torsion to the free pole, a rush of electric tertiary induction is heard, whose value is  $40^\circ$ . Again, testing this wire by moving the coil, I find only a remaining magnetism of 10, and upon repeating the experiment of elastic torsion, I find a tertiary of 5; thus we can go on gradually discharging the wire, but it will be found that its discharge is a recombination, and that it first passed through the stage I have mentioned.

Heat has a very great effect upon molecular magnetic effects. On iron it increases the current, but in steel the current is diminished. For experimenting on iron wire, which gave a tertiary current of  $50^\circ$  positive (with a torsion of  $20^\circ$ ), upon the application of the flame of a spirit-lamp, the force rapidly increases (care being taken not to approach red heat), until the force is doubled, or  $100^\circ$  positive. The same effects were obtained in either direction, and were not due to a molar twist or thermo-current, as if care had been taken to put on not more than  $10^\circ$  of torsion, the wire came back to zero at once on removal of the torsion. Hard tempered steel, whose value was  $10^\circ$  whilst cold, with a torsion of  $45^\circ$ , became only  $1^\circ$  when heated, but returned (if not too much heated) to  $8^\circ$  when cold. I very much doubted this experiment at first, but on repeating the experiment with steel several times, I found that on heating it, I had softened the extreme hard (yellow) temper to that of the well-known blue temper. Now, at blue temper, hot, the value of steel was but  $1^\circ$  to  $2^\circ$ , whilst soft iron of a similar size gave  $50^\circ$  of force cold, and  $100^\circ$  at red heat. Now, as I have already shown that the effects I have described depend on molecular elasticity, it proves at least, as far as iron and steel are concerned, that a comparatively perfect elastic body, such as tempered steel, has but slight molecular elasticity, and that heat reduces it, but that soft iron, having but little molar elasticity, has a molecular elasticity of a very high degree, which is increased by heat.

The objects of the present paper being to bring the experimental facts before the notice of the Royal Society, and not to give a theoretical solution of the phenomena, I will simply add, that if we assume with Poisson, that the paths of the molecules of iron are circles, and that they become ellipses by compression or strain, and also that they are capable of being polarised, it would sufficiently explain the new effects.

Joule has shown that an iron bar is longer and narrower during magnetisation than before, and in the case of the transverse strain, the exterior portions of the wire are under a far greater strain than those

near the centre, and as the polarised ellipses are at an angle with the molecules of the central portions of the wire, its polarisation reacts upon them, producing the comparatively strong electric currents I have described.

### III. "On the Action of Sodium upon Chinoline." By C. GREVILLE WILLIAMS, F.R.S. Received March 8, 1881.

In 1867 I made some experiments on the action of sodium upon chinoline and lepidine, and found that a substance was produced which dyed silk a beautiful but fugitive orange colour. I announced this fact in a paper "On the Higher Homologues of Chinoline."<sup>\*</sup> I made analyses of the products at the time, but the difficulties in the way of preparing them pure were so great, and the time at my disposal was so limited, that I did not make public any quantitative results until March, 1878, when I published a short note "On the Action of Sodium on Chinoline and Lepidine."<sup>†</sup> In that paper I gave the results of an analysis of a product from chinoline, which agreed with the formula  $C^{18}H^{14}N^2 \cdot HCl$ , which is obviously that of the hydrochlorate of dichinoline. I also gave an analysis of the nitrate of dilepidine, which agrees almost perfectly with the theoretical values; but I did not enter into the details of the modes of preparation. In the course of my investigation of  $\beta$ -lutidine, it was natural that I should study the action of sodium upon it, but I met with so many and unexpected difficulties that I determined to prepare myself for a new attack upon the subject by a fresh investigation of the action of sodium upon chinoline. As I find that other observers are working upon chinoline,  $\beta$ -lutidine, and  $\beta$ -collidine,<sup>‡</sup> I have thought it desirable to bring before the Society the results obtained, although the investigation is still proceeding.

#### *Action of Sodium upon Chinoline.*

The action of sodium on chinoline is exceedingly remarkable, not merely because it polymerises the base, for a similar result, as is well known, takes place with picoline, but because the products have properties which are, I think, different from any yet observed among organic substances. For a yellow oil like dichinoline to yield a true although fugitive dye, in the form of a brilliantly red crystalline

\* "Laboratory," May, 1867, p. 109.

† "Chemical News," March 1, 1878.

‡ Richard, "Bull. Soc. Chim.," [2] xxxii, 486—489; Boutlerow and Wischnegradsky, *loc. cit.*, No. 9, June 5, 1880; Oechsner de Coninck, "Bull. Soc. Chim.," Nos. 4 and 5, p. 210, September 5, 1880. The latter chemist has repeated several of my older experiments, evidently under the impression that they had not been made before.